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The environmental dependence of gas accretion on to galaxies: quenching satellites through starvation

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ABSTRACT

Galaxies that have fallen into massive haloes may no longer be able to accrete gas from their surroundings: a process referred to as ‘starvation’ or ‘strangulation’ of satellites. We study the environmental dependence of gas accretion on to galaxies using the cosmological, hydrodynamical EAGLE simulation. We quantify the dependence of gas accretion on stellar mass, redshift, and environment, using halo mass and galaxy overdensity as environmental indicators. We find a strong suppression, of many orders of magnitude, of the gas accretion rate in dense environments, primarily for satellite galaxies. This suppression becomes stronger at lower redshift. However, the scatter in accretion rates is very large for satellites. This is (at least in part) due to the variation in the halocentric radius, since gas accretion is more suppressed at smaller radii. Central galaxies are influenced less strongly by their environment and exhibit less scatter in their gas accretion rates. The star formation rates of both centrals and satellites show similar behaviour to their gas accretion rates. The relatively small differences between gas accretion and star formation rates demonstrate that galaxies generally exhaust their gas reservoir somewhat faster at higher stellar mass, lower redshift, and in denser environments. We conclude that the environmental suppression of gas accretion could directly result in the quenching of star formation.

Key words: methods: numerical – galaxies: evolution – galaxies: formation – galaxies: haloes – intergalactic medium – galaxies: star formation.

1 INTRODUCTION

Star-forming galaxies are observed to lie on a fairly tight relation between their star formation rate (SFR) and stellar mass, and there appears to be an equilibrium between gas supply and gas removal (e.g. Noeske et al. 2007; Bouché et al. 2010; Schaye et al. 2010; Lee et al. 2015). Quenched galaxies are gas-poor and have SFRs significantly below this ‘star-forming main sequence’. The mechanisms proposed to be responsible for such low gas masses and SFRs include the removal of the galaxies’ interstellar medium (ISM) as well as the prevention of further gas accretion on to the galaxies.

The ISM can be removed through feedback from star formation or supermassive black holes or through environmental effects, such as ram pressure, tidal stripping or satellite–satellite encounters. When the ISM is removed efficiently, this results in a fast decline of the SFR. The removal of a galaxy’s gaseous halo through ram pressure can result in the suppression of gas accretion on to the galaxy and a slow decline of the SFR, also known as ‘starvation’ or ‘strangulation’ (e.g. Larson, Tinsley & Caldwell 1980; Balogh, Navarro & Morris 2000). All of these mechanisms are thought to play a role in the quenching of galaxies, but which one dominates likely depends on the stellar mass and environment (e.g. van den Bosch et al. 2008; Fabello et al. 2012; Bahé & McCarthy 2015). In order to keep galaxies quenched for an extended period, resulting in the observed red colours, their gas accretion needs to be halted or significantly reduced compared to that of blue, star-forming galaxies. This may

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be caused by a different process than the one that quenches galaxies in the first place or there could be multiple quenching mechanisms acting together (e.g. Bahé & McCarthy 2015; Marasco et al. 2016; Trayford et al. 2016).

Observationally, galaxy properties such as colour, morphology and specific SFR have been shown to depend on stellar mass (e.g. Kauffmann et al. 2003; Baldry et al. 2004; Peng et al. 2010; Moustakas et al. 2013). Similarly, there is a clear dependence of these galaxy properties on the environment, at least at low redshift (e.g. Dressler 1980; Kauffmann et al. 2004; Peng et al. 2010; Wetzel, Tinker & Conroy 2012). Whether or not similar environmental dependences exist at $z > 1$, where the cosmic SFR density is significantly higher, is still an open question, but most studies find weaker environmental trends (e.g. Scoville et al. 2013; Lin et al. 2016).

The environment of a galaxy is thought to be primarily determined by properties of the dark matter-dominated halo it lives in (e.g. Crain et al. 2009) and commonly used environmental indicators all correlate strongly with the halo mass (Haas, Schaye & Jeason-Daniel 2012; Marasco et al. 2016). Indeed, using observed galaxy group catalogues, it has been shown that colour and star formation history correlate mostly with the mass of the host halo (e.g. Blanton & Berlind 2007; Wilman, Zibetti & Budavári 2010). The virial radius, R_{vir} , of a halo approximately marks the location of an accretion shock, within which the shocked gas is in virial equilibrium with the dark matter. For sufficiently massive systems, the gas temperature in the halo is much higher than that in the unshocked intergalactic medium. The temperature can be further increased through galactic feedback (e.g. van de Voort & Schaye 2012; Fielding et al. 2016; van de Voort et al. 2016). Central galaxies sit at the minimum of the host halo’s potential well, where the gas densities and cooling rates are highest, whereas satellite galaxies have fallen into more massive haloes and are thus offset from the centre of the potential.

The quenching mechanisms, ram pressure stripping and starvation, are special in the sense that they only operate on satellite galaxies. When a galaxy falls into a massive halo, it can lose its own gaseous halo and thus lose the reservoir from which it can replenish its ISM (e.g. Larson et al. 1980; Balogh et al. 2000; Peng, Maiolino & Cochrane 2015). Starvation therefore results in relatively slow quenching as the galaxy slowly consumes its cold gas reservoir through star formation. When the pressure of the external hot gas is high enough to also remove the ISM of the satellite galaxy, this is referred to as ram pressure stripping (e.g. Gunn & Gott 1972; Quilis, Moore & Bower 2000; Roediger & Brüggén 2007; Tonnesen & Bryan 2009; Steinhauser, Schindler & Springel 2016) and results in relatively fast quenching of star formation.

Gas accretion, or the lack thereof, is difficult to observe directly, but it affects the SFR, which is a key observable for galaxies. Additionally, a decrease in the infall of pristine gas can explain the fact that the metallicity is observed to increase with decreasing SFR at fixed stellar mass (Mannucci et al. 2010; Davé, Oppenheimer & Finlator 2011; Peng et al. 2015; Kacprzak et al. 2016; Lagos et al. 2016; Bahé et al. 2017). To directly study gas accretion, cosmological, hydrodynamical simulations are a powerful tool and they enable us to quantify the effect of starvation in different environments. Simha et al. (2009) indeed found lower gas accretion rates for satellites in massive haloes compared to centrals of the same mass. The current generation of state-of-the-art cosmological simulations has both the resolution and the range of environments required to study the dependence of gas accretion on halo mass, stellar mass and galaxy overdensity in detail. This provides an important means

to understand the quenching of galaxies in general and of satellites in particular.

In this paper, we use the Evolution and Assembly of GaLaxies and their Environments (EAGLE) simulation (Crain et al. 2015; Schaye et al. 2015) to study the gas accretion rates of central and satellite galaxies and how they depend on stellar mass, environment and redshift. Relatedly, the removal of cold gas through ram pressure and tidal stripping in the same simulation was studied by Marasco et al. (2016), who showed that EAGLE reproduces the observed H I mass–environment trends.

In Section 2, we describe the simulation used and the way we define galaxies and haloes (Section 2.1) and gas accretion on to galaxies (Section 2.2). In Section 3, we present our results. We start by using halo mass as the environmental indicator, showing results for central and satellite galaxies separately in Section 3.1. Section 3.2 shows how accretion rates vary with halocentric distances. We then adopt a geometric definition for environment (as often used in observations) in Section 3.3, namely galaxy overdensity, based on the 10th nearest neighbour, which confirms our previous results, albeit with small modifications. In Section 3.4, we show the implications for the resulting SFRs. The evolution with redshift is discussed throughout Section 3. Finally, we discuss and summarize our results in Section 4.

2 METHOD

This work is part of the EAGLE project (Schaye et al. 2015; Crain et al. 2015), which consists of a large number of cosmological simulations with varying subgrid physics, size, and resolution. The simulations use an extensively modified version of GADGET-3 (last described in Springel 2005), a smoothed particle hydrodynamics (SPH) code. The main modifications are the formulation of SPH, the time stepping and the subgrid physics. For our main results, we make use of the largest-volume simulation (*‘Ref-L100N1504’*). The model is comprehensively described in Schaye et al. (2015) and we will summarize its main properties here.

The simulations assume a Λ CDM cosmology with parameters $\Omega_m = 1 - \Omega_\Lambda = 0.307$, $\Omega_b = 0.04825$, $h = 0.6777$, $\sigma_8 = 0.8288$, and $n = 0.9611$ (Planck Collaboration I et al. 2014). A cubic volume with periodic boundary conditions is defined, within which the mass is distributed over 1504^3 dark matter and an equal number of gas particles. The volume of the simulation used in this work is $(100 \text{ comoving Mpc})^3$. The (initial) particle masses for baryons and dark matter are $1.8 \times 10^6 \text{ M}_\odot$ and $9.7 \times 10^6 \text{ M}_\odot$, respectively. The gravitational softening length is $2.66 \text{ comoving kpc}$ but is limited to a maximum value of 0.7 proper kpc , which is reached at $z = 2.8$.

For gas with hydrogen number densities $n_H \geq 10^{-5} \text{ cm}^{-3}$, a temperature floor of 8000 K is imposed. Additionally, a temperature floor T_{EoS} is imposed at all densities. This takes the form of a polytropic equation of state $P_{\text{tot}} \propto \rho_{\text{gas}}^{4/3}$, where P_{tot} is the total pressure and ρ_{gas} is the density of the gas, normalized to $T_{\text{EoS}} = 8000 \text{ K}$ at $n_H = 0.1 \text{ cm}^{-3}$.

In dense gas, star formation is modelled according to the recipe of Schaye & Dalla Vecchia (2008). The SFR per unit mass depends on the gas pressure and reproduces the observed Kennicutt–Schmidt law (Kennicutt 1998) by construction. Following Dalla Vecchia & Schaye (2012), a gas is eligible to form stars when it has a temperature within 0.5 dex of T_{EoS} and a density exceeding $n_H^* = 0.1 \text{ cm}^{-3} (\frac{Z}{0.002})^{-0.64}$, where Z is the gas metallicity. Because the true, multiphase structure of the ISM is not resolved, the gas temperature reflects the effective pressure and the density should be interpreted as the mean density of the ISM. In this work, the galaxies’ ISM is

therefore defined to be all of the star-forming gas (i.e. all particles with $\text{SFR} > 0$).

The abundances of 11 elements released by massive stars and intermediate mass stars are followed as described in Wiersma et al. (2009b). We assume the stellar initial mass function of Chabrier (2003), ranging from 0.1 to $100 M_{\odot}$. Radiative cooling and heating are computed element-by-element in the presence of the cosmic microwave background radiation and the Haardt & Madau (2001) model for the UV/X-ray background from galaxies and quasars, assuming the gas to be optically thin and in (photo)ionization equilibrium, as in Wiersma, Schaye & Smith (2009a).

Feedback from star formation is implemented using the prescription of Dalla Vecchia & Schaye (2012) by injecting energy in thermal form. Neighbouring gas particles are selected stochastically (depending on the available energy) and then heated by a fixed temperature increment of $10^{7.5}$ K. The amount of energy available depends on the local gas density and metallicity, and its parameters were chosen to approximately reproduce the observed $z = 0$ galaxy stellar mass function and galaxy sizes (Crain et al. 2015; Schaye et al. 2015).

The model for the formation of and feedback from supermassive black holes is fully described in Rosas-Guevara et al. (2015) and Schaye et al. (2015). A seed mass black hole of about $1.4 \times 10^5 M_{\odot}$ is placed in every halo with mass $M_{\text{halo}} \gtrsim 1.4 \times 10^{10} M_{\odot}$ that does not already contain a black hole (Springel, Di Matteo & Hernquist 2005). These black holes grow by the accretion of gas, after which energy is injected into the surrounding medium, and by mergers. A fraction of 1.5 per cent of the rest-mass energy of the gas accreted on to the black hole is injected into the surrounding medium in the form of heat, by increasing the temperature of at least one neighbouring gas particle by at least $10^{8.5}$ K. The parameter values were chosen to obtain a good match with the observed galaxy masses and sizes as well as with the $z = 0$ relation between galaxy mass and black hole mass (Crain et al. 2015).

These simulations were not tuned to reproduce any observable gas properties. However, they have been shown to reproduce a number of observations reasonably well, such as the H I column density distribution and covering fraction (Rahmati et al. 2015), the H I and H₂ contents of galaxies with stellar mass above $10^{10} M_{\odot}$ (Lagos et al. 2015; Bahé et al. 2016; Crain et al. 2017) and the dependence of the H I mass on environment (Marasco et al. 2016). However, the predicted soft X-ray luminosity and halo gas mass fraction in groups and clusters are somewhat higher than observed (Schaye et al. 2015), whilst there may not be sufficient cool gas in low-mass galaxies (Crain et al. 2017). The simulation also reproduces the observed colour–magnitude diagram (Trayford et al. 2015), the FIR properties of galaxies (Camps et al. 2016), their optical colours (Trayford et al. submitted), alpha enhancements (Segers et al. 2016b) and the evolution of the galaxy mass function (Furlong et al. 2015) and galaxy sizes (Furlong et al. 2017).

2.1 Identifying haloes and galaxies

In order to identify haloes, we first find dark matter haloes using a Friends-of-Friends (FoF) algorithm. If the separation between two dark matter particles is less than 20 per cent of the average dark matter particle separation, they are placed in the same group. Baryonic particles are placed in a group if their nearest dark matter neighbour is part of the group. We then apply SUBFIND (Springel et al. 2001; Dolag et al. 2009) to the FoF output to find the gravitationally bound particles and to identify subhaloes. The galaxy containing the minimum of the gravitational potential is called the central. Note that

when the central and satellite are of very similar mass, there is some ambiguity in the assignment of either label.

In this paper, we use the most bound particle of a halo or subhalo, as found by SUBFIND, as the halo or galaxy centre. We then use a spherical overdensity criterion, considering all the particles in the simulation, to determine the halo mass, M_{halo} , of every host halo (thus excluding subhaloes). Here, we define M_{halo} to be the total mass enclosed by a radius, R_{vir} , within which the mean overdensity is 200 times the mean density of the Universe at its redshift. Satellites are subhaloes of a larger structure identified by SUBFIND and may or may not have their own dark matter halo. We include only those that reside within the R_{vir} of their host halo in order to exclude spurious satellite identifications, but our results are insensitive to this choice. Stellar mass, M_{star} , is measured using the star particles associated with the galaxy’s subhalo, but only including those within 30 proper kpc of the galaxy’s centre in order to exclude intrahalo stars. For each galaxy, we calculate the ISM mass by adding up the mass of star-forming gas particles in the galaxy’s subhalo, again within 30 kpc of the centre.

2.2 Calculating gas accretion on to galaxies

Gas accretion at $z_1 > z > z_2$ is calculated in the following way. We select the gas particles that were not part of any galaxy’s ISM at z_1 , i.e. those that were not star forming. This automatically excludes accretion of ISM gas through galaxy mergers, but still allows accretion of non-star-forming gas previously surrounding external galaxies. Of those selected particles, we consider the gas and recently formed star particles that are in the galaxy of interest at z_2 (and were non-star-forming gas at z_1) as having been accreted between z_1 and z_2 . Gas particles that accrete after z_1 but are ejected or stripped before z_2 are therefore not included in our accretion rates. This affects only the normalization of our results, but not the trends. When approximately doubling the time interval between z_1 and z_2 , we find that the normalization of the gas accretion rates changes by 0.1 dex or less. We therefore conclude that the contribution from accreted but rapidly ejected gas is likely subdominant. We include all accreted gas, regardless of its history, which therefore includes accretion from an outer non-star-forming gas disc as well as direct accretion from the gaseous halo. The gas accretion rate differs from the change in gas mass even in the absence of star formation and stellar mass-loss, because the removal of gas that is in the ISM at z_1 , through stripping or outflows, is not taken into account.

For brevity, accretion at $0.1 > z > 0$ ($2.24 > z > 2$) will be referred to as accretion at $z = 0$ ($z = 2$). This redshift interval corresponds to 1.35 Gyr (0.33 Gyr) in our assumed cosmology. The minimum non-zero accretion rate we can measure for an individual galaxy is therefore $1.3 \times 10^6 M_{\odot} \text{ Gyr}^{-1}$ at $z = 0$ and $5.5 \times 10^6 M_{\odot} \text{ Gyr}^{-1}$ at $z = 2$ (i.e. one accreted gas particle over the time interval). Halo mass or host mass refers to the mass of the host halo in which the satellite resides at z_2 , unless otherwise stated. We repeated our analysis at $1.26 > z > 1$ and found it to show behaviour intermediate between $z = 0$ and $z = 2$, although more similar to $z = 2$ and therefore chose not to include it in this paper.

3 RESULTS

As described in the introduction, gas accretion may proceed very differently for satellite and central galaxies. Satellites may be stripped of their own haloes and no longer able to replenish their ISM through accretion. In this section, we investigate the importance of this starvation and its dependence on stellar mass and redshift, as

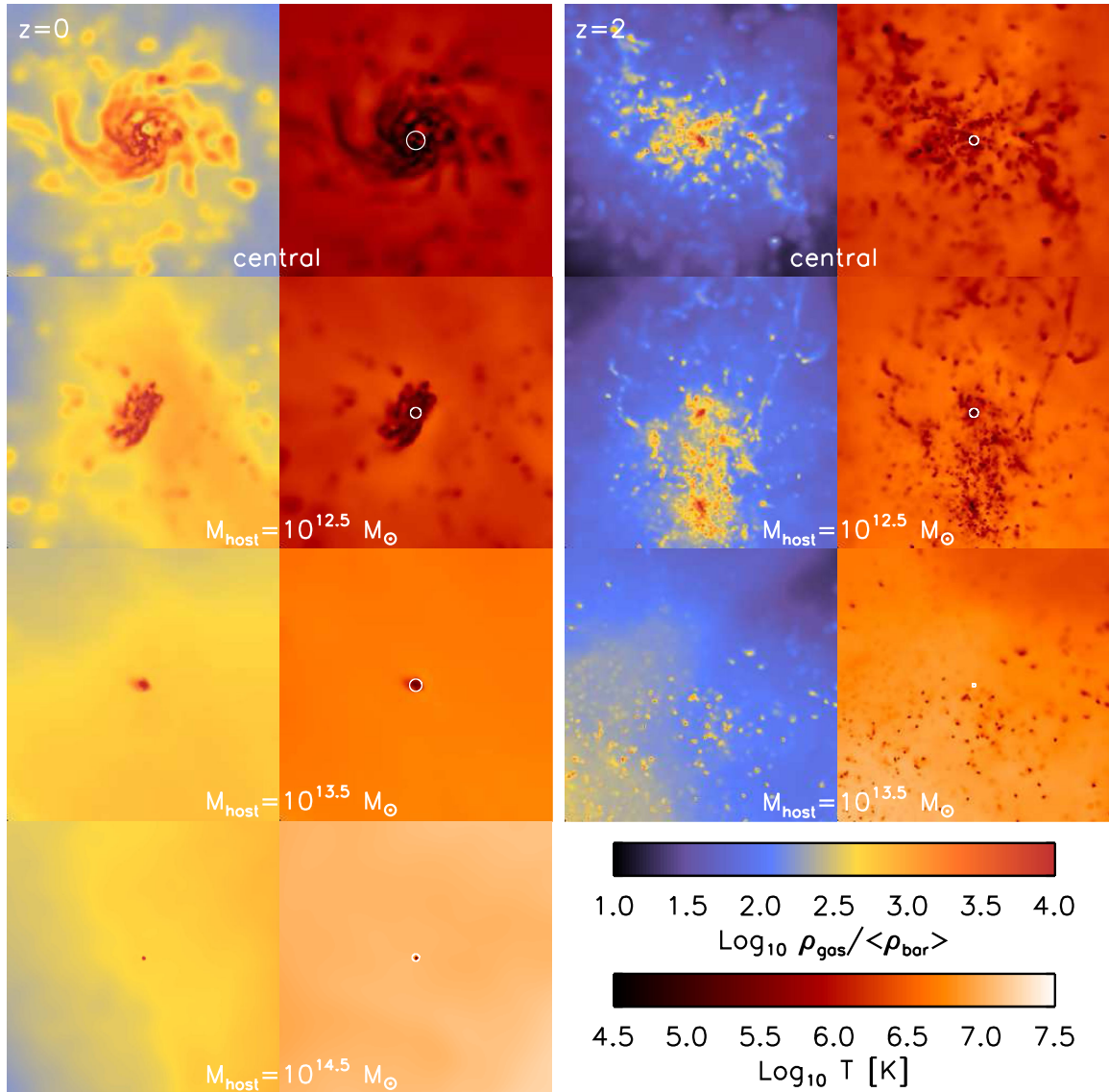


Figure 1. Projected $(200 \text{ proper kpc})^3$ images of the average gas overdensity (first and third columns) and temperature (second and fourth columns) at $z = 0$ (left) and $z = 2$ (right), each centred on a randomly chosen galaxy with $M_{\text{star}} = 10^{10} M_{\odot}$. White circles in the temperature panels indicate the galaxies' stellar half-mass radii. The top row shows central galaxies, while the other rows show satellite galaxies for which the host halo mass increases from $10^{12.5} M_{\odot}$ (second row) to $10^{13.5} M_{\odot}$ (third row) to $10^{14.5} M_{\odot}$ (bottom row). The temperature of the surrounding gas increases with host halo mass, which increases the effect of the environment on gas accretion rates. At high redshift, there is much more cold, dense gas around the galaxies. Gas at the same overdensity also has a factor of 27 higher density at $z = 2$ than at $z = 0$, because of the cosmological expansion of the Universe. Together, these effects result in a decrease of the environmental dependence with increasing redshift.

well as on environment, which we separate into a dependence on halo mass (Sections 3.1 and 3.4), halocentric radius (Section 3.2), and galaxy overdensity (Section 3.3).

Regarding gas accretion on to galaxies, an important difference between high and low redshift is the density of the circumgalactic medium, which controls the cooling rate, and thus affects the amount of cold, dense gas. At fixed redshift, the gas accretion rate is affected by the environment through differences in the temperature and pressure of the hot, diffuse gas and the amount of cold, dense gas around a galaxy. Fig. 1 shows these gas properties visually. Each image is a projection of the same physical volume, $(200 \text{ proper kpc})^3$. The galaxy in the centre of each image has approximately the same stellar mass, $M_{\text{star}} = 10^{10} M_{\odot}$, and its stellar half-mass radius (computed within 30 proper kpc) is indicated by

a white circle. The first and third columns show volume-weighted average gas overdensity, and the second and fourth columns show mass-weighted average gas temperature, at $z = 0$ (left) and $z = 2$ (right). The images in the top row show central galaxies and the other three rows show satellite galaxies for which the host halo mass increases from $M_{\text{host}} = 10^{12.5} M_{\odot}$ (second row) to $10^{13.5} M_{\odot}$ (third row) to $10^{14.5} M_{\odot}$ (bottom row, only at $z = 0$). The satellites' (proper) halocentric distances are 186 kpc ($10^{12.5} M_{\odot}$), 318 kpc ($10^{13.5} M_{\odot}$), and 710 kpc ($10^{14.5} M_{\odot}$) at $z = 0$ and 108 kpc ($10^{12.5} M_{\odot}$) and 205 kpc ($10^{13.5} M_{\odot}$) at $z = 2$.

We selected these galaxies randomly, but verified that others show very similar gas properties (except that the mean density and, to a lesser extent, the mean temperature depend on the distance from the centre of the halo). The gaseous environments of central

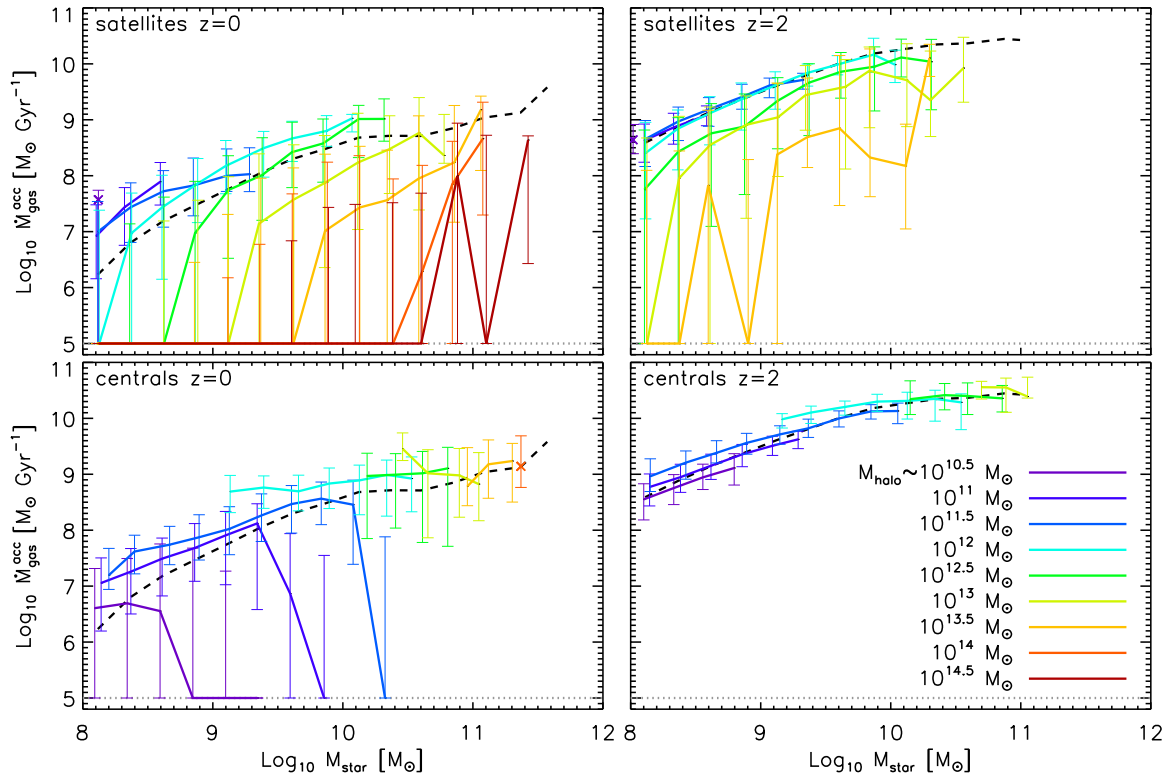


Figure 2. Median gas accretion rate on to galaxies as a function of stellar mass for satellite (top) and central galaxies (bottom) at $z = 0$ (left) and $z = 2$ (right). The dashed curves show median accretion rates for all galaxies. Solid curves and crosses show median accretion rates in halo mass bins, 0.5 dex in width, centred on the values indicated in the legend. Gas accretion rates of zero have been set to $10^5 \text{ M}_{\odot} \text{Gyr}^{-1}$, as shown by the dotted, grey line. Error bars show the 16th and 84th percentiles of the distribution. Only bins containing at least five galaxies are shown. Gas accretion on to satellites is more strongly suppressed in more massive host haloes and for lower mass galaxies. Although much weaker, this trend is reversed for central galaxies. The environmental dependence becomes stronger and the scatter becomes larger at lower redshift.

galaxies and satellites in $10^{12.5} \text{ M}_{\odot}$ host haloes are similar (first and second row in Fig. 1), so $M_{\text{star}} = 10^{10} \text{ M}_{\odot}$ satellites may not be affected much by their environment. However, the temperature of the surrounding gas increases strongly with host halo mass (third and fourth row), which also means that the average pressure and the average cooling time increase (at fixed density). This makes it harder for a satellite to maintain its own gaseous halo and to accrete cold gas on to its ISM.

As expected, there is a large difference between high and low redshift, with a lot more cold, dense gas surrounding the galaxies at $z = 2$ in all environments. These denser satellite haloes will be less easily stripped. Note that there is a factor of 27 difference in physical density between gas at fixed overdensity at $z = 2$ and $z = 0$, because the average density of the universe decreases with decreasing redshift. At higher redshift, a larger fraction of the accreting gas never shock heats to the virial temperature near the virial radius of the halo, known as ‘cold-mode’ accretion (e.g. Kereš et al. 2005; Dekel et al. 2009; van de Voort et al. 2011; Correa et al. in preparation). Average (gas and dark matter) densities are higher and cooling times are shorter, which results in higher gas accretion rates on to a galaxy. The availability of gas with short cooling times likely reduces the effect of the (hot halo) environment.

3.1 Dependence on halo mass

The environmental effects on gas accretion, discussed qualitatively in the previous section, are shown quantitatively in Fig. 2. The

median gas accretion rate on to galaxies is shown as a function of stellar mass for satellite (top) and central galaxies (bottom) at $z = 0$ (left) and $z = 2$ (right). The dashed curves show median accretion rates for all galaxies and are thus identical in the top and bottom panels. Solid curves show median accretion rates in halo mass bins, 0.5 dex in width, centred on the values indicated in the legend and increasing from purple ($10^{10.5} \text{ M}_{\odot}$) to red ($10^{14.5} \text{ M}_{\odot}$). This halo mass is the mass of the host halo in which the satellite resides; it is not the mass of the satellite’s subhalo. For visualization purposes, gas accretion rates of zero have been set to $10^5 \text{ M}_{\odot} \text{Gyr}^{-1}$ (as also indicated by the grey, dotted line). Note that we cannot detect accretion of less than 1 particle per redshift interval, so the ‘zero’ rates could, in reality, be higher ($< 1.3 \times 10^6 \text{ M}_{\odot} \text{Gyr}^{-1}$ at $z = 0$ and $< 5.5 \times 10^6 \text{ M}_{\odot} \text{Gyr}^{-1}$ at $z = 2$). Thin, coloured error bars show the 1σ scatter (i.e. 16th and 84th percentiles of the distribution). Only bins containing at least five galaxies are shown.

In general, the gas accretion rate increases with the stellar mass, a consequence of the increase of the dark matter and gas accretion rate on to haloes (e.g. Neistein, van den Bosch & Dekel 2006; Fakhouri & Ma 2010; van de Voort et al. 2011; Correa et al. 2015). It increases more strongly with the stellar mass at the low-mass end than at the high-mass end. This can be explained by the increase of the cooling time of the circumgalactic medium at higher masses (with higher virial temperatures) and by the mass dependence of stellar and active galactic nucleus (AGN) feedback (e.g. van de Voort et al. 2011), which sets the shape of the galaxy stellar mass function.

Satellites of fixed stellar mass accrete gas at rates that differ by many orders of magnitude, depending on the mass of their host halo. At fixed stellar mass, the suppression of gas accretion becomes gradually stronger with increasing halo mass. This is due to the fact that the temperature and pressure of the halo gas as well as the velocity of the satellite increase with mass, which makes the removal of the satellite’s gaseous subhalo more efficient and the accretion of gas on to the satellite more difficult. Conversely, at fixed halo mass, the suppression of gas accretion becomes weaker with increasing stellar mass. This is due to the fact that the temperature and pressure differences between the satellite’s subhalo and its host halo are smaller and that the potential well of the satellite is deeper, which makes stripping of the satellite’s subhalo less efficient and gas accretion easier. Lower mass satellites have also been shown to spend more time in their host halo before they merge or disrupt due to the reduced dynamical friction. This means that low-mass satellites at a given redshift have on average spent more time as satellites, increasing the likelihood of their subhaloes being stripped away (e.g. Boylan-Kolchin, Ma & Quataert 2008; Wetzel & White 2010). Similar trends with stellar mass and environment are seen for the H I content of satellites (Marasco et al. 2016).

On the other hand, central galaxies behave very differently. The dependence on halo mass at fixed stellar mass is much weaker and the gas accretion rate is slightly higher in more massive haloes (opposite to the trend for satellites). Since the central resides at the minimum of the potential, gas cooling and accretion can happen efficiently and increase with halo mass. The higher gas accretion rate will lead to a higher SFR and thus, eventually, to a higher stellar mass. In this way, central galaxies regulate their own growth, which keeps the scatter in the ‘star-forming main sequence’ and the stellar mass–halo mass relation small. In general, the accretion rate increases with stellar mass, but the median accretion rate is very low for the more massive galaxies in the three lowest halo mass bins ($M_{\text{halo}} \lesssim 10^{11.5} M_{\odot}$). These haloes’ gas reservoirs are too small to efficiently feed the central galaxy. Additionally, these galaxies have formed more stars than lower mass galaxies in similar haloes and thus produced stronger outflows, which prevents accretion at least temporarily.

The scatter in gas accretion rates is very large for galaxies that are strongly influenced by their environment. This results (at least in part) from the fact that some satellites have fallen into the host halo recently and have not yet been affected, whereas others have spent considerably more time inside the halo and were stripped of their own halo (and possibly their ISM). This time dependence will be studied by Bahé et al. (in preparation). This means, however, that just knowing the stellar mass and host halo mass is not enough to predict a galaxy’s accretion rate. Observationally, galaxies in cluster environments show a very large range of SFRs at fixed stellar mass (Chies-Santos et al. 2015; Paccagnella et al. 2016), which potentially results from the large range in gas accretion rates for satellites.

As expected, gas accretion rates on to galaxies are higher at $z = 2$ than at $z = 0$, following increased dark matter and gas accretion rates on to haloes (e.g. Neistein et al. 2006; Fakhouri & Ma 2010; van de Voort et al. 2011; Correa et al. 2015). The dependence on halo mass is weaker at $z = 2$, because gas densities in the circumgalactic medium are higher at high redshift and the accretion of cold, dense gas is less readily affected by the environment than that of hot, diffuse gas. Additionally, dynamical friction is stronger, which means that galaxies merge more quickly at higher redshift and thus spend less time as satellites. None the less, even at $z = 2$, gas accretion on to satellites is substantially suppressed at $M_{\text{halo}} \gtrsim 10^{12.5} M_{\odot}$. The

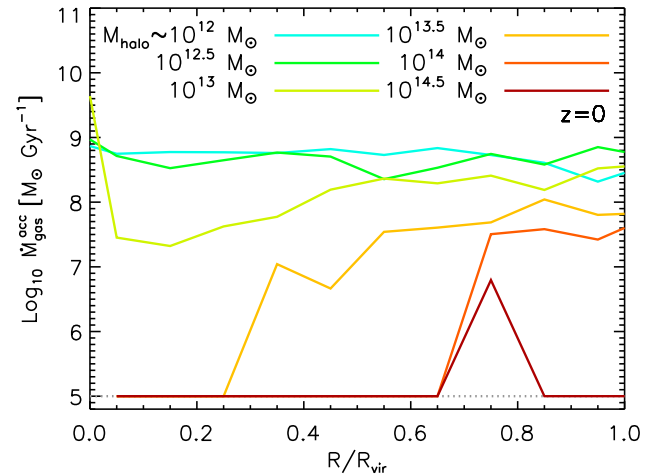


Figure 3. Median $z = 0$ gas accretion rate on to galaxies with $10^{9.5} < M_{\text{star}} < 10^{10.5} M_{\odot}$ as a function of 3D distance from the host halo centre, normalized by their host halo’s virial radius. Different coloured curves correspond to different halo mass bins (0.5 dex wide) as indicated in the legend. Gas accretion rates of zero have been set to $10^5 M_{\odot} \text{ Gyr}^{-1}$, as shown by the dotted, grey line. Bins contain at least five galaxies and only central galaxies are included in the first radial bin. For high-mass haloes, the gas accretion on to galaxies decreases by orders of magnitude towards the centre of the host halo. This may indicate that environmental effects are stronger at smaller radii or that galaxies near the centre have spent more time in their host halo or both.

trends with environment are similar to those at $z = 0$, except for the absence of massive centrals in low-mass host haloes with low accretion rates. Observations show a reduced quenching of satellite galaxies towards higher redshift, which also indicates a smaller environmental effect (e.g. Hatfield & Jarvis 2016; Kavinwanichakij et al. 2016; Nantais et al. 2016).

3.2 Radial dependence

It is expected that the suppression of gas accretion due to the host halo environment and thus the quenching of star formation are slow processes (e.g. McCarthy et al. 2008; Simha et al. 2009; Wetzel et al. 2013). On average, satellites in the outskirts of the halo have likely fallen in more recently than those in the centre of the host halo (Gao et al. 2004). This should then result in a radial dependence of the gas accretion rates with respect to the halo centre. Fig. 3 shows the $z = 0$ median gas accretion rate for galaxies with $10^{9.5} < M_{\text{star}} < 10^{10.5} M_{\odot}$ as a function of the 3D distance from the host halo centre, normalized by the virial radius of the halo. The environmental dependence is shown by different curves, which include haloes within 0.25 dex of the values indicated in the legend. For clarity, median gas accretion rates of zero are shown at $10^5 M_{\odot} \text{ Gyr}^{-1}$. Note that this calculation uses 3D radii while any observed radial dependence would be diluted by projection effects. The radial trends are similar for galaxies with lower (higher) stellar masses than $\approx 10^{10} M_{\odot}$, but the suppression of the gas accretion rate at fixed halo mass and radius is stronger (weaker), as expected from Fig. 2. Similar, but weaker, radial trends also exist at higher redshift.

The two lowest halo mass bins that contain $M_{\text{star}} \approx 10^{10} M_{\odot}$ satellite galaxies (cyan and green curves; $M_{\text{halo}} \approx 10^{12-12.5} M_{\odot}$) show no radial dependence of the gas accretion rates on to satellite galaxies. This is consistent with Fig. 2, where we found little suppression of gas accretion for these masses. However, a strong radial

dependence can be seen for $M_{\text{halo}} \gtrsim 10^{13} M_{\odot}$, with the median accretion rate decreasing towards smaller radii. At fixed (normalized as well as physical) radius, the accretion rate is lower in more massive haloes and the radius where the median accretion rate reaches zero in our simulation moves to higher values. This radial dependence may indicate that environmental effects are stronger at smaller halocentric radii or that galaxies near the centre have spent more time in their host halo or both. Knowing the position with respect to the centre of the halo can thus help reduce the scatter in gas accretion rate at fixed stellar and halo masses.

In our most massive haloes ($M_{\text{halo}} \gtrsim 10^{13.5} M_{\odot}$), the accretion rate of these satellites is, on average, suppressed throughout the entire halo and it is possible that the halo's influence extends beyond its virial radius. The dependence outside R_{vir} will be studied in Bahé et al. (in preparation).

If starvation is a major contributor to satellite quenching, this radial dependence of the gas accretion rate will result in a similar radial dependence of the SFR. Paccagnella et al. (2016) indeed found that, at fixed stellar mass, the SFRs are lower in cluster centres than in their outskirts (see also Wetzel et al. 2012). Although this is consistent with our results, other environmental quenching mechanisms, such as ram pressure stripping or tidal stripping, would lead to similar behaviour.

3.3 Dependence on galaxy overdensity

Halo mass can be inferred in observations through various methods, such as galaxy clustering, weak lensing, satellite kinematics, and abundance matching. However, most of these methods are indirect and may introduce biases. To facilitate comparison to observations, we use the galaxy's 10th nearest neighbour to define environment in this section. This quantity is more directly available in galaxy surveys, although we measure the 3D galaxy density, whereas observations will necessarily use the projected 2D density and a relative velocity cut along the line of sight. Our resulting galaxy density, ρ_{10} , is defined as the number density of galaxies above our resolution limit (i.e. $M_{\text{star}} > 10^8 M_{\odot}$). The results do not change substantially if instead we choose to use e.g. the 5th nearest neighbour or $M_{\text{star}} > 10^9 M_{\odot}$. We then compute the average galaxy density in the full simulation volume, ρ_{all} , and define the galaxy overdensity as $\delta \equiv \rho_{10}/\rho_{\text{all}}$ and use this as our environmental parameter.

Fig. 4 is similar to Fig. 2 but uses δ instead of M_{halo} to characterize a galaxy's environment. It shows median gas accretion rates as a function of stellar mass for satellites (top), centrals (middle) and all galaxies (bottom) at $z = 0$ (left) and $z = 2$ (right). In all panels, the dashed curves show median gas accretion rates for all galaxies in our simulation. Solid curves show median gas accretion rates in bins of galaxy overdensity, as indicated in the legend and error bars show the 1σ scatter.

The conclusions drawn from this definition of environment, i.e. galaxy overdensity, are very similar to those drawn above based on the halo mass. This is expected, because galaxy overdensity correlates with halo mass (e.g. Haas et al. 2012; Marasco et al. 2016). Considering all galaxies first (bottom panels), the gas accretion rate is lower in denser environments at fixed stellar mass. At fixed galaxy overdensity, the gas accretion rate is more suppressed for lower stellar masses. Note, however, that the scatter in accretion rate can be very large, likely due in part to the varying amounts of time a galaxy spent in a particular environment. The top panels show that this environmental effect is primarily driven by satellites, as also seen in observations of the environmental dependence of

the quenched galaxy fraction (Peng et al. 2012). We find that the environmental effect decreases with increasing redshift (see also Section 3.1).

Satellites are affected by their environment more than centrals, but the effect on centrals is non-zero for large overdensities, $\delta \gtrsim 30$, at $z = 0$. We did not see this gas accretion suppression for centrals when we used halo mass as a probe of environment. Although the similarity with Fig. 2 shows that halo mass is responsible for the main environmental effect in satellite galaxies, this implies a (smaller) residual environmental dependence besides halo mass. This may be due to the fact that, in dense environments, nearby galaxies are competing for the same gas reservoir and accrete gas that would otherwise have been available to other galaxies (Aragon-Calvo, Neyrinck & Silk 2014). Relatedly, there may be environmental effects beyond R_{vir} , either because the halo affects the intergalactic medium beyond R_{vir} or because a satellite's orbit brings it beyond R_{vir} rendering it a central once more (i.e. splashback galaxies). Additionally, galaxies in dense environments formed earlier, also known as assembly bias (e.g. Sheth & Tormen 2004; Gao, Springel & White 2005), which reduces their growth rate at later times.

The environmental dependence of the gas accretion rate for central galaxies in our simulation is possibly reversed at the highest stellar masses at $z = 2$, although this is a very small effect. Centrals at a fixed stellar mass $> 10^{10.5} M_{\odot}$ have slightly higher gas accretion rates in higher density environments. This reversal of the trend could lead to a reversal in the SFR–density relation, which has been suggested observationally at $z = 1 - 1.5$ (Elbaz et al. 2007; Lin et al. 2016), but only if the galaxies are predominantly centrals, which is likely the case for massive galaxies.

3.4 Effect on star formation

SFRs are primarily determined by the amount of gas in the ISM, which is replenished through gas accretion, gas-rich galaxy mergers, and stellar mass-loss. It is difficult to measure gas accretion rates observationally, but SFRs are routinely derived. Because gas accretion and star formation are so intricately linked, we turn our attention to the balance between them and the dependence on environment.

The environmental dependence of star formation (and ISM mass) is very similar to the environmental dependence of gas accretion. At fixed environment (either halo mass or galaxy overdensity), the gas accretion rate tends to be somewhat more suppressed than the SFR. This is due to the fact that there is some time delay during which galaxies consume their existing ISM, before they are affected by starvation. Moreover, stellar mass-loss can replenish the ISM even in the absence of accretion (Leitner & Kravtsov 2011; Segers et al. 2016a).

Instead of showing the environmental effect on star formation directly, which is qualitatively the same as the effect on gas accretion shown in previous figures, Fig. 5 shows the ratio of the median gas accretion rate and median SFR as a function of stellar mass for satellites (top) and centrals (bottom) at $z = 0$ (left) and $z = 2$ (right). The SFR is calculated as the stellar mass in a galaxy that was formed between $z = 0.1$ ($z = 2.24$) and $z = 0$ ($z = 2$) divided by the corresponding time interval. The dashed curves show gas-accretion-rate-to-SFR ratios for all galaxies. Solid curves show these ratios in halo mass bins, 0.5 dex in width, centred on the values indicated in the legend. We show only stellar mass bins where the median gas accretion rate is non-zero, because we cannot resolve low-level gas accretion and this would otherwise result in biased and unrealistic

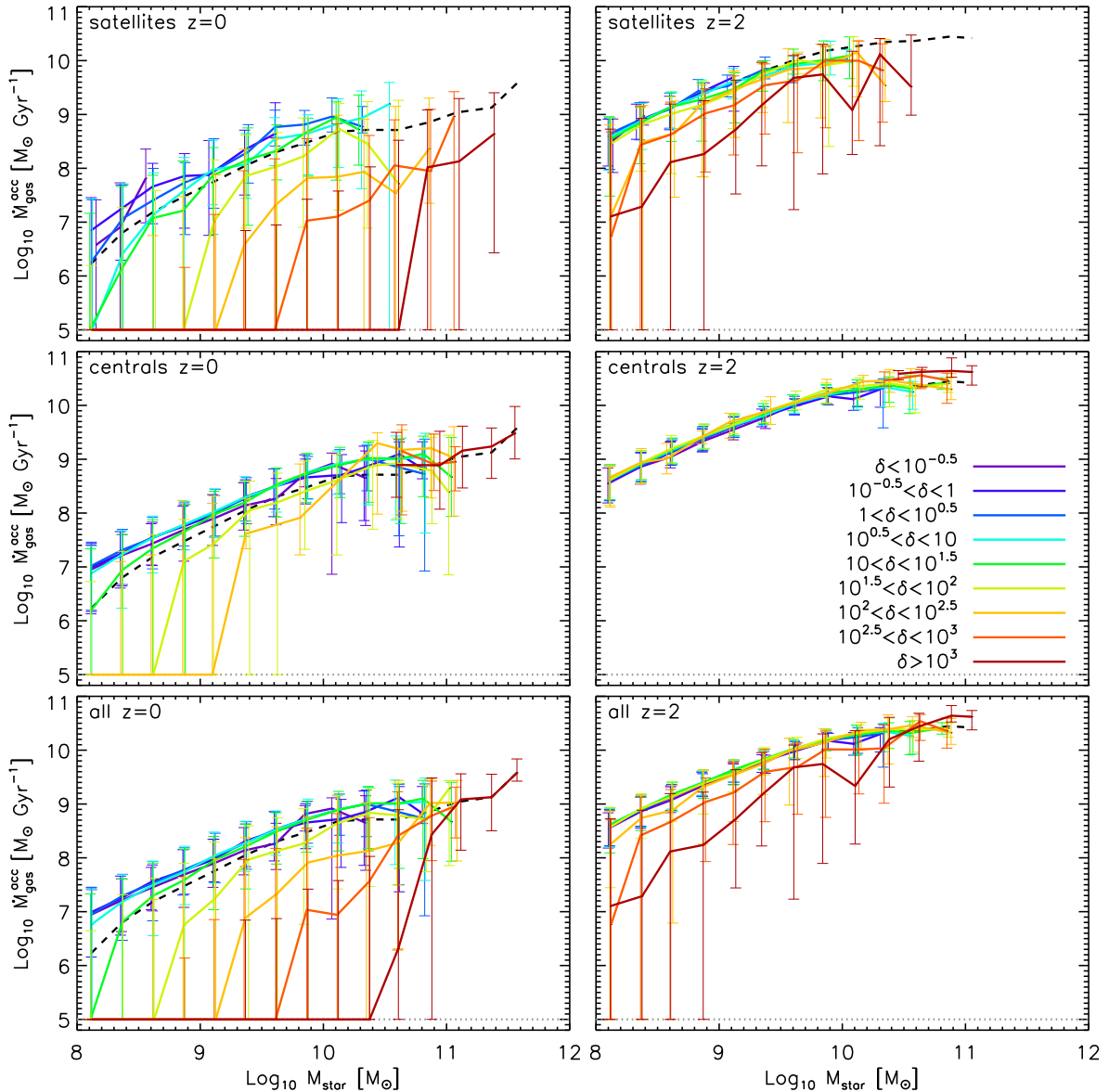


Figure 4. Median gas accretion rates as a function of stellar mass for satellites (top), centrals (middle) and all galaxies (bottom) at $z = 0$ (left) and $z = 2$ (right). Dotted and dashed curves are the same as in Fig. 2. Solid curves show median gas accretion rates in bins of galaxy overdensity, $\delta \equiv \rho_{10}/\rho_{\text{all}}$, as indicated in the legend. Gas accretion rates of zero have been set to $10^5 \text{ M}_{\odot} \text{ Gyr}^{-1}$. Error bars show the 16th and 84th percentiles of the distribution. Only bins containing at least five galaxies are shown. There is a very strong dependence on galaxy overdensity at $z = 0$ and a weaker dependence at $z = 2$. This dependence is primarily driven by satellite galaxies, although there is some residual dependence for centrals at $z = 0$.

or undefined ratios. Although this excludes the lowest-stellar mass bins at fixed halo mass, no galaxies have been excluded for the stellar mass bins shown and our results are therefore not biased.

The gas-accretion-rate-to-SFR ratio generally decreases with increasing stellar mass. At $z = 2$, the most massive galaxies have ratios close to unity, on average, which implies that their ISM mass stays roughly constant in the absence of outflows, stripping and mergers. Lower mass galaxies, however, still have gas accretion rates higher than their SFRs and are building up their gas reservoirs. Low-mass galaxies at $z = 0$ also have gas accretion rates that approximately balance their SFRs. Most $M_{\text{star}} > 10^{10} \text{ M}_{\odot}$ galaxies at $z = 0$ are depleting their ISM. This is consistent with the downsizing trend that is found in observations and simulations, where more massive galaxies assemble their stellar mass earlier than less massive ones (e.g. Bundy et al. 2006; Pérez-González et al. 2008; Crain

et al. 2009; Muzzin et al. 2013; Okamoto, Shimizu & Yoshida 2014; van de Voort 2016). At fixed halo mass, more massive centrals have lower gas-accretion-rate-to-SFR ratios. However, satellites in massive haloes ($M_{\text{halo}} \gtrsim 10^{13} \text{ M}_{\odot}$) at $z = 0$ exhaust their ISM less quickly at higher stellar mass, because the effect of the environment on gas accretion and star formation decreases with increasing stellar mass. It is important to note that this comparison includes only gas accretion and star formation and excludes gas brought in by galaxy mergers and gas removed through outflows or stripping. The change in gas mass may therefore differ from the change implied by the gas-accretion-rate-to-SFR ratio.

The weaker dependence on stellar mass and halo mass (up to 1 dex) as compared to the gas accretion rate (Fig. 2) shows that the rate of gas accretion on to galaxies is one of the primary drivers of the SFR. Stripping of the ISM may also contribute or even dominate

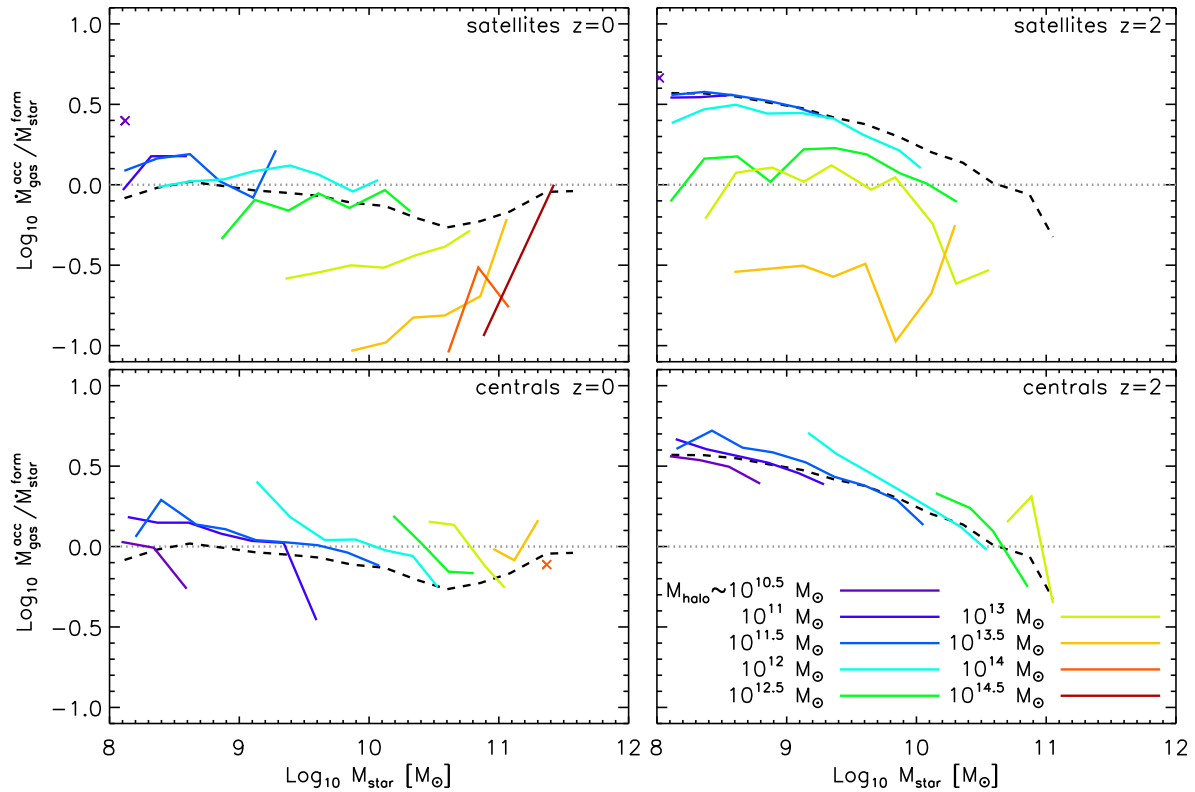


Figure 5. The median gas accretion rate divided by the median SFR as a function of stellar mass for satellites (top) and central galaxies (bottom) at $z = 0$ (left) and $z = 2$ (right). The dashed curves show the gas-accretion-rate-to-SFR ratios for all galaxies. Solid curves and crosses show the gas-accretion-rate-to-SFR ratio in halo mass bins, 0.5 dex in width, centred on the values indicated in the legend. Only bins containing at least five galaxies and non-zero median gas accretion rates are shown. The dotted line shows a ratio of unity, where the median gas accretion rate and the median SFR are equal. There is only a relatively weak dependence of this ratio on stellar mass and environment, showing that the SFR is, in large part, driven by the gas accretion rate. In general, the ratio decreases with stellar mass, which means that massive galaxies exhaust their gas reservoirs more quickly than dwarf galaxies at $z = 0$ and build up their gas reservoir more slowly at $z = 2$. In denser environments (in this case in more massive host haloes), satellite galaxies have lower gas-accretion-rate-to-SFR ratios, by up to 1 dex, whereas centrals have slightly higher ratios.

in certain regimes. Even though gas accretion rates and SFRs show similar environmental dependences, the gas-accretion-rate-to-SFR ratio clearly shows a residual dependence on environment. At fixed stellar mass, satellites have lower gas-accretion-rate-to-SFR ratios if they reside in higher mass haloes, indicating that their gas reservoir is building up more slowly ($M_{\text{halo}} \lesssim 10^{13} M_{\odot}$ at $z = 2$) or is being depleted more rapidly ($M_{\text{halo}} \gtrsim 10^{13} M_{\odot}$ at $z = 2$ and $M_{\text{halo}} \gtrsim 10^{12.5} M_{\odot}$ at $z = 0$). This is likely because of the time delay between the shut-down of gas accretion and that of star formation due to the remaining gas reservoir in the ISM and the replenishment by stellar mass-loss. Centrals, on the other hand, exhibit a smaller environmental dependence of opposite sign. At fixed stellar mass, centrals in more massive haloes have slightly higher gas-accretion-rate-to-SFR ratios, which allows them to grow their ISM more rapidly, which (after some time delay) results in higher SFRs and higher stellar masses, keeping the scatter in the stellar mass–halo mass relation for central galaxies small (Matthee et al. 2017; Tinker et al. 2016).

4 DISCUSSION AND CONCLUSIONS

We have used the EAGLE simulation to study how the rate of gas accretion on to galaxies changes with environment. A galaxy’s environment is defined either as the mass of the host halo in which it resides or as the local 3D galaxy overdensity based on its 10th

nearest neighbour with $M_{\text{star}} > 10^8 M_{\odot}$. The most significant environmental effect is seen for satellites, whose gas accretion rates vary by many orders of magnitude (see Figs 2 and 4). At fixed stellar mass, the suppression of gas accretion is larger in more massive haloes and at higher galaxy overdensity, albeit with large scatter. This can be caused by the more extreme gas properties in higher mass haloes, such as higher temperature and pressure and fewer cold, dense clumps, by the higher satellite velocities and/or by the fact that satellites spend a larger amount of time in these harsher environments before they merge or disrupt. We also find a strong dependence on the halocentric distance for satellites, which have lower gas accretion rates at smaller radii (see Fig. 3). Similarly, this is likely also caused by the higher gas pressure in the halo centre, the higher velocities and the longer time spent inside the halo. This radial dependence can also help explain the large scatter seen for satellites with similar stellar and host halo masses.

The effect of the environment decreases towards higher redshift, both at fixed halo mass and at fixed galaxy overdensity. Since the gas density is higher at higher redshift, the cooling time of the gas is shorter, and a larger fraction of the gas is cold and dense. Satellite galaxies are therefore less easily stripped of their own gaseous halo and can more easily accrete gas from their surroundings. Additionally, dynamical friction is stronger and galaxies thus merge more quickly at higher redshift, which causes them to spend less time as satellites, reducing environmental effects.

Although satellites are affected more strongly by their environment, centrals are affected as well. At fixed stellar mass, centrals in more massive haloes have slightly higher gas accretion rates, opposite to the trend for satellites. Since more massive haloes have larger gas reservoirs, this results in higher accretion rates in their centres and also acts to keep the scatter in the halo-to-stellar mass relation low for central galaxies. Remarkably, for higher galaxy overdensities, we find a suppression of the gas accretion rate for centrals at $z = 0$, similar to (but less strong than) the effect of overdensity on satellites. This means that there is some environmental dependence not determined by the halo mass, possibly because strongly clustered central galaxies formed earlier compete for the same gas reservoir and/or were satellites previously.

Stellar and AGN feedback can reduce the gas accretion rate on to galaxies by an order of magnitude (Faucher-Giguère, Kereš & Ma 2011; van de Voort et al. 2011). This is due to the fact that outflows do not only eject ISM gas from galaxies, but also plough into inflowing gas, thus preventing it from accreting. Although feedback is necessarily modelled with simplified subgrid recipes in EAGLE, the fact that this simulation reproduces the galaxy stellar mass function lends credence to the idea that the implemented subgrid models capture the large-scale effects of the outflows reasonably well. Particularly relevant for our study is that Marasco et al. (2016) have shown that EAGLE reproduces the observed dependence of $H\text{I}$ content on environment. Even if the effect of feedback is not captured well and galactic winds behave differently in reality, this will primarily affect the normalization of the relations discussed in this work, but the trends with stellar mass, environment and redshift are likely robust to changes in feedback models.

The suppression of gas accretion in dense environments has direct bearing on our understanding of the quenching of satellite galaxies if they are indeed (in part) quenched through starvation. There is some observational evidence for long quenching times, which could point to starvation as the quenching mechanism, since galaxies will slowly consume their remaining ISM until they run out of fuel for star formation (e.g. Wetzel et al. 2013; Wheeler et al. 2014; Peng et al. 2015). The large scatter in gas accretion rates at fixed stellar mass can potentially explain the large range in observed SFRs in cluster environments (e.g. Chies-Santos et al. 2015; Paccagnella et al. 2016). The smaller effect of the environment on galaxy SFRs observed at higher redshift (e.g. Scoville et al. 2013; Hatfield & Jarvis 2016; Kawinwanichakij et al. 2016; Lin et al. 2016; Nantais et al. 2016) is consistent with the reduced environmental dependence of the gas accretion rates found in our simulation. Since, for a fixed environment, gas accretion is more strongly suppressed in lower mass galaxies, we expect that quenching of their star formation proceeds more rapidly, if starvation indeed plays an important role in galaxy quenching.

The SFR in our simulation follows the gas accretion rate closely. It exhibits the same trends with environment, although it is somewhat less suppressed than the gas accretion rate. This is likely due to the fact that a galaxy is not immediately stripped of its ISM and is replenished by stellar mass-loss, so that it can continue to form stars even if its gas supply is completely shut off until it runs out of fuel. This leads us to conclude that, at least for small galaxies in dense environments, the quenching of star formation could directly result from the suppression of gas accretion, which is strongly affected by the properties of the surrounding gas. However, we did not investigate the removal of the ISM through environmental processes, such as ram pressure stripping or tidal stripping, which could dominate the quenching of some or all satellites. In the absence of stripping, outflows and mergers, most galaxies at $z = 2$ are building

up their ISM, whereas massive galaxies at $z = 0$ are depleting it. More massive galaxies and satellites in denser environments generally deplete their gas reservoirs more quickly, consistent with the observed downsizing of galaxies and the SFR–density relation (see Fig. 5). However, at fixed halo mass, satellites in groups and clusters at $z = 0$ deplete their ISM less quickly with increasing stellar mass due to the decreasing environmental dependence of their gas accretion and star formation rates.

The environmental dependences found in our hydrodynamical simulation could help inform and improve the treatment of gas accretion on to satellites in semi-analytic models (e.g. Guo et al. 2016). We focused on the ‘instantaneous’ gas accretion rates for central and satellite galaxies. As hinted at by the dependence on halocentric radius (see Fig. 3), the accretion rate on to a satellite is affected by the amount of time spent in the host halo and/or its trajectory. The time-scales over which the gas accretion rates decrease after galaxies fall into massive haloes will be explored in future work (Bahé et al. in preparation), which will help to understand the origin of the large scatter in accretion rates shown in this paper.

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APPENDIX A: RESOLUTION TEST

For our resolution test, we use a simulation with the same resolution as the one used for our main results, but with a volume that is smaller by a factor of 64 (*‘Ref-L25N376’*). The simulation volume limits the total number of galaxies and which environments we can probe, but there is no volume effect at fixed environment. We have repeated our analysis at 8 (2) times higher mass (spatial) resolution using a simulation with the same feedback parameters (*‘Ref-L25N752’*) and one with recalibrated parameters to again match the stellar mass function and galaxy sizes (*‘Recal-L25N752’*). Note that the former simulations do not match these observables as well for numerical reasons (Schaye et al. 2015). We find that the normalization of the gas accretion rates depends on resolution. At fixed resolution, it depends on the (recalibrated) feedback parameters, but only for massive galaxies ($M_{\text{star}} \gtrsim 10^{10} M_{\odot}$). Higher resolution leads to increased gas densities, which result in higher cooling rates and higher accretion rates. The difference in gas accretion rates between our standard-resolution simulation and the higher resolution simulations is 0.3 dex on average.

Despite this normalization issue, the trends with stellar mass, halo mass and galaxy overdensity remain the same, as shown in Fig. A1. This figure shows the median gas accretion rate as a function of stellar mass for simulation *‘Ref-L25N376’* (left) and *‘Ref-L25N752’* (right) at $z = 0$. Satellites (centrals) are shown in the top (bottom)

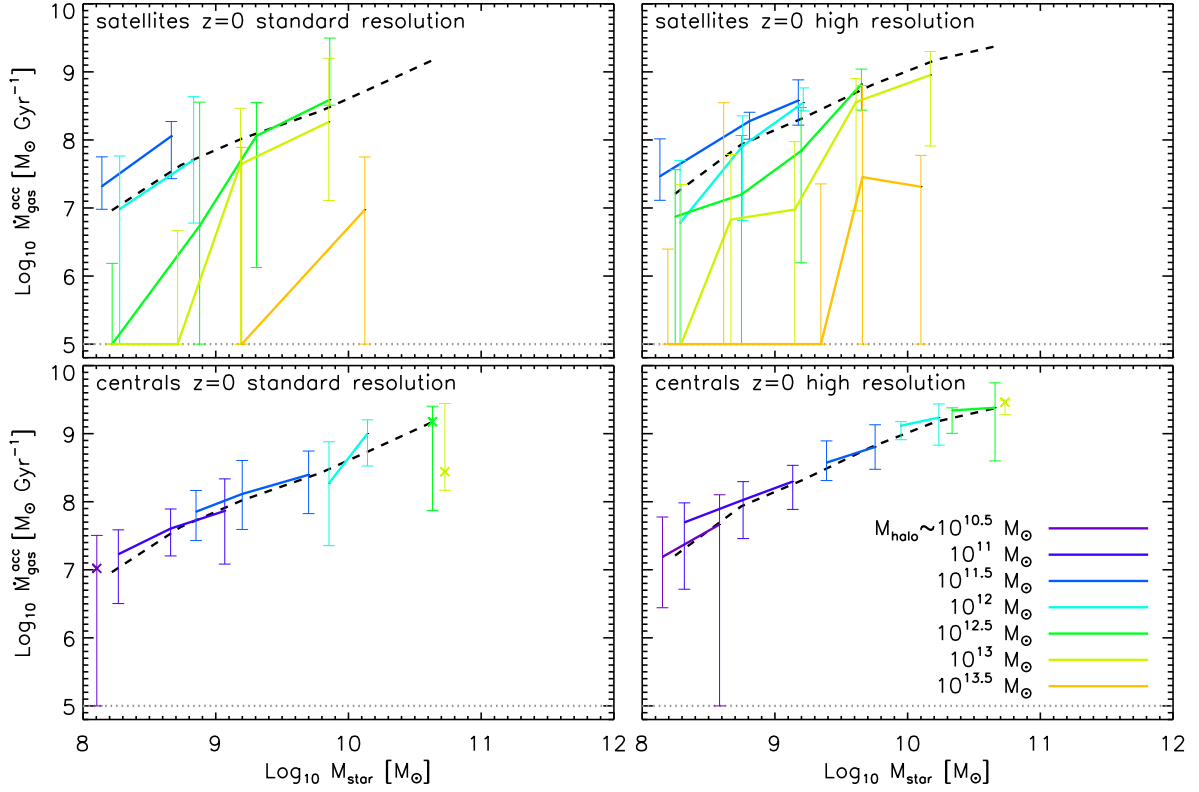


Figure A1. Median gas accretion rate on to galaxies as a function of stellar mass for a $(25 \text{ comoving Mpc})^3$ simulation at standard resolution (*'Ref-L25N376'*; left) and a $(25 \text{ comoving Mpc})^3$ simulation at eight times higher mass resolution (*'Ref-L25N752'*; right) for satellites (top) and central galaxies (bottom) at $z = 0$. The dashed curves show median accretion rates for all galaxies. Solid, coloured curves and crosses show median accretion rates in halo mass bins, 0.5 dex in width, centred on the values indicated in the legend. Gas accretion rates of zero have been set to $10^5 \text{ M}_\odot \text{ Gyr}^{-1}$. Error bars show the 16th and 84th percentiles of the distribution. The bins are twice as large in M_{star} as in Fig. 2, to increase statistics, and only bins containing at least five galaxies are shown. The suppression of gas accretion on to satellites due to the environment, visible in the top panels, is independent of resolution.

panels in different halo-mass bins, as shown by the coloured curves. The dashed curves show median accretion rates for all galaxies in each simulation. Gas accretion rates of zero have been set to $10^5 \text{ M}_\odot \text{ Gyr}^{-1}$ and the error bars show the scatter around the median. The environmental suppression of gas accretion in satellites is insensitive to resolution. We therefore conclude that the behaviour

of the gas accretion on to galaxies, described in this paper, is robust. Note, however, that we are unable to probe cluster environments in these small simulation volumes.

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